

Polytrope & White dwarfs

$$K = c_1 A^{-2} G \rho_c^{\frac{n-1}{n}} \quad (\text{KH 19.9}),$$

c_i : non dimensional constants

$$\& \rho_c \propto \bar{\rho} \propto MR^{-3} = M \left(\frac{A}{z_n} \right)^3$$

$$\rightarrow A = c_2 \rho_c^{1/3} M^{-1/3}$$

$$\rightarrow K = c_3 M^{2/3} G \rho_c^{1-\frac{1}{n}-\frac{2}{3}} = c_3 M^{2/3} G \rho_c^{\frac{1}{3}-\frac{1}{n}},$$

$$\text{also: } P_c = K \rho_c^{1+\frac{1}{n}} = c_3 G M^{2/3} G \rho_c^{\frac{4}{3}}$$

$$P_c = 0.48 GM^{2/3} G \rho_c^{\frac{4}{3}} \quad (n = \frac{3}{2}) \quad (1.21)$$

$$= 0.36 GM^{2/3} G \rho_c^{\frac{4}{3}} \quad (n = 3) \quad (1.22)$$

Lane – Emden equation

$$\frac{1}{z^2} \frac{d}{dz} \left(z^2 \frac{dw}{dz} \right) + w^n = 0$$

$$M = 4\pi \rho_c R^3 \left(-\frac{1}{z} \frac{dw}{dz} \right)_{z=z_n}$$

$$\frac{\bar{\rho}}{\rho_c} = \left(-\frac{3}{z} \frac{dw}{dz} \right)_{z=z_n}$$

n	z_n	$\left(-z^2 \frac{dw}{dz} \right)_{z=z_n}$	$\bar{\rho}/\bar{\rho}$
0	2.4494	4.8988	1.0000
1	3.14159	3.14159	3.28987
1.5	3.65375	2.71406	5.99071
2	4.35287	2.41105	11.40254
3	6.89685	2.01824	54.1825
4	14.97155	1.79723	622.408
4.5	31.8365	1.73780	6189.47
5	∞	1.73205	∞

準静的収縮

Quasi-static Contraction

$$E_{GR} = -\int_0^R \frac{Gm\rho}{r} 4\pi r^2 dr = -\frac{3GM^2}{5R} \quad (\rho = \text{一定の時})$$

$$\text{一般に} \quad E_{GR} \approx -\frac{GM^2}{R}$$

$$(1.5) \rightarrow P_c \ll P \gg \frac{GM^2}{4\pi R^4} \quad (1.8)$$

$$\rho_c \text{は星の平均密度に比例} \quad \rho_c \propto \frac{M}{\frac{4}{3}\pi R^3}$$

$$\frac{P_c}{\rho_c} \propto \frac{GM}{R} \quad (1.9), \quad \frac{P_c^3}{\rho_c^4} \propto G^3 M^2 \quad (1.10)$$

$$\text{理想気体状態方程式} \quad P_c = \frac{k}{\mu m_u} \rho_c T_c$$

$$+(1.10) \rightarrow \frac{T_c^3}{\rho_c} \propto G^3 M^2 \quad (1.11)$$

$$(1.11) \text{式} \Rightarrow M \text{一定のとき} \quad \rho_c \propto T_c^3$$

Quasi-static contraction

エントロピー(単位体積あたりの)

$$s = \frac{k}{\mu m_u} \ln(T^{3/2} / \rho) + C_1,$$

+ (1.11) \rightarrow

$$\left(\frac{\mu m_u}{k} \right) s_c = \ln\left(\frac{M^2}{T_c^{3/2}} \right) + C_2 = \ln\left(\frac{M}{\rho_c^{1/2}} \right) + C_3 \quad (1.12)$$

(1.12)式 $\Rightarrow M$ 一定で s_c が減少すると T_c 上昇

$$\text{比熱は負} \quad C_s = \frac{ds_c}{d \ln T_c} = -1.5 \frac{k}{\mu m_u} \quad (1.13)$$

Main Sequence

• Maximum Stellar temperature determined by the degenerate pressure (rough estimate)

(1.6) ビリアル定理 $2E_{KE} + E_{GR} = 0$

$$E_{GR} \approx -\frac{GM^2}{R}, \quad E_{KE} = \frac{3}{2}NkT \rightarrow kT \approx \frac{GMm_u}{3R} \approx Gm_u M^{2/3} \rho^{1/3}$$

+ 縮退圧の効く密度: $\rho \sim \frac{M}{R^3} \sim \rho$ (プロイ波長)

$$\lambda \approx \frac{h}{\sqrt{m_e kT}}, \quad \rho \approx \frac{m_u}{\lambda^3} \approx m_u \frac{(m_e kT)^{3/2}}{h^3}$$

$$kT_{max} \approx \left[\frac{G^2 m_u^{8/3} m_e}{h^2} \right] M^{4/3} \quad (1.14)$$

Hydrogen Burning (KH193)

- (1.14) $\rightarrow M < 0.08 M_\odot - T_{max} < \sim 10^7 \text{ K}$
Hydrogen can't be ignited.
- \rightarrow Brown dwarfs (褐色矮星, supported by the degenerate pressure)
- $T \sim 1.5 \times 10^7 \text{ K}$: **pp (proton-proton) chain**

$$p + p \rightarrow d + e^+ + \nu_e$$

$$p + d \rightarrow {}^3\text{He} + \gamma$$

- ${}^3\text{He} + {}^3\text{He} \rightarrow {}^4\text{He} + p + p$ (pp1)
- ${}^3\text{He} + {}^4\text{He} \rightarrow {}^7\text{Be} + \gamma$

- $e^- + {}^7\text{Be} \rightarrow {}^7\text{Li} + \nu_e$

$$p + {}^7\text{Li} \rightarrow {}^4\text{He} + {}^4\text{He}$$
 (pp2)

- $p + {}^7\text{Be} \rightarrow {}^8\text{B} + \gamma$

$${}^8\text{B} \rightarrow {}^8\text{Be}^* + e^+ + \nu_e$$

$${}^8\text{Be}^* \rightarrow {}^4\text{He} + {}^4\text{He}$$
 (pp3)
$$\epsilon_{pp} \sim 2.36 \times 10^6 X_H^2 \rho T_6^{-2/3} \text{ exp}(-33.8 T_6^{-1/3}) \text{ erg g}^{-1} \text{ sec}^{-1}$$

Solar neutrino (KH338)

- pp (proton-proton) chain**

$$p + p \rightarrow d + e^+ + \nu_e$$

$$p + d \rightarrow {}^3\text{He} + \gamma$$

- ${}^3\text{He} + {}^3\text{He} \rightarrow {}^4\text{He} + p + p$ (pp1)
- ${}^3\text{He} + {}^4\text{He} \rightarrow {}^7\text{Be} + \gamma$

- $e^- + {}^7\text{Be} \rightarrow {}^7\text{Li} + \nu_e$

$$p + {}^7\text{Li} \rightarrow {}^4\text{He} + {}^4\text{He}$$
 (pp2)

- $p + {}^7\text{Be} \rightarrow {}^8\text{B} + \gamma$

$${}^8\text{B} \rightarrow {}^8\text{Be}^* + e^+ + \nu_e$$

$${}^8\text{Be}^* \rightarrow {}^4\text{He} + {}^4\text{He}$$
 (pp3)

Process	Neutrino flux ($10^{14} \text{ m}^{-2} \text{ s}^{-1}$)	Maximum neutrino energy (MeV)
$p + p \rightarrow d + e^+ + \nu_e$	$6.0 (1 \pm 0.02)$	0.420
$e^- + {}^7\text{Be} \rightarrow {}^7\text{Li} + \nu_e$	$0.47 (1 \pm 0.15)$	0.861
${}^8\text{B} \rightarrow {}^8\text{Be}^* + e^+ + \nu_e$	$5.8 \times 10^{-4} (1 \pm 0.37)$	15
${}^{13}\text{N} \rightarrow {}^{13}\text{C} + e^+ + \nu_e$	$0.06 (1 \pm 0.50)$	1.199
${}^{15}\text{O} \rightarrow {}^{15}\text{N} + e^+ + \nu_e$	$0.05 (1 \pm 0.58)$	1.732

$\nu_e + {}^{37}\text{Cl} \rightarrow {}^{37}\text{Ar} + e^+$ (detectable only above 0.81 MeV)
Predicted $7.9 \pm 2.6 \text{ SNU}$, Observed $2.5 \pm 0.4 \text{ SNU}$

Kamioka $\nu_e + {}^2\text{H} \rightarrow p + p + e^-$, $\nu_x + {}^2\text{H} \rightarrow n + p + \nu_x$
Result ν_e changes to ν_{mu} & ν_{tau}

Discovery of the neutrino oscillation (& neutrino mass)

Hydrogen burning

- $M > 1 M_\odot$, at $T \sim 2 \times 10^7 \text{ K}$
- CNO cycle (KH196)**

$$\epsilon_{CNO} \sim 8.67 \times 10^{27} X_{CNO} X_H \rho T_6^{-2/3} \text{ exp}(-152.3 T_6^{-1/3}) \text{ erg g}^{-1} \text{ sec}^{-1}$$

$$\propto \rho T^{16} \text{ at } T \sim 2 \times 10^7 \text{ K}$$

${}^{12}\text{C} : {}^{13}\text{C} : {}^{14}\text{N} : {}^{15}\text{N} = 4 : 1 : 95 : 0.004$

\Rightarrow heavier elements can be produced a little